## d. Example 4

Obtain an expression for the entropy change in an RK gas when the gas is isothermally compressed. Determine the entropy change when superheated R–12 is isothermally compressed at 60°C from 0.0194 m<sup>3</sup> kg<sup>-1</sup> (state 1) to 0.0126 m<sup>3</sup> kg<sup>-1</sup> (state 2). Compare the result with the tabulated value of  $s_1 = 0.7259, \, s_2 = 0.6881$ .

Solution

Consider the RK state equation

$$P = RT/(v-b) - a/(T^{1/2}v(v+b))$$
(A)

Note that the attractive force constant a is different from "a" Helmholtz function. From the third relation Eq. (22) and Eq. (A),

$$(\partial s/\partial v)_T = (\partial P/\partial T)_v = R/(v-b) + (1/2) a/(T^{3/2}v(v+b)).$$
 (B)

Integrating Eq. (B),

$$s_2(T,v_2) - s_1(T,v_1) =$$

$$R\ln((v_2-b)/(v_1-b)) + (1/2)(a/(T^{3/2}b)) \ln(v_2(v_1+b)/(v_1(v_2+b))).$$
 (C)

From Table 1 for R–12,  $T_c = 385$  K, and  $P_c = 41.2$  bar. Therefore  $\overline{a}$  =208.59 bar (m<sup>3</sup> kmole<sup>-1</sup>) <sup>2</sup> K<sup>1/2</sup>, and  $\overline{b}$  = 0.06731 m<sup>3</sup> kmole<sup>-1</sup>. The molecular weight M = 120.92 kg kmole<sup>-1</sup>, and

$$\begin{aligned} \mathbf{a} &= \overline{\mathbf{a}}/M^2 = 208.59 \text{ bar } (m^3 \text{ kmole}^{-1}) \text{ }^2K^{1/2} \div 120.92 \text{ }^2(\text{kg kmole}^{-1})^2 \\ &= 1.427 \text{ k Pa } (m^3 \text{ kg}^{-1}) \text{ }^2K^{1/2}, \text{ and} \\ \mathbf{b} &= \overline{\mathbf{b}}/M = 0.557 \times 10^{-3} \text{ m}^3 \text{ kg}^{-1}. \\ \text{Since, } R &= 8.314 \div 120.92 = 0.06876 \text{ kJ kg}^{-1} \text{ K}^{-1}, \\ \mathbf{s}_2 - \mathbf{s}_1 &= 0.06876 \ln[(0.0126 - 0.000557) \div (0.0194 - 0.000557)] \\ &+ (1/2) \{172.5 \div (333^{1.5} \ 0.000557)\} \ln [0.0126 \ (0.0194 + 0.000557) \\ &\div \{0.0194 \times (0.0126 + 0.000557)\}]. \\ &= -0.06876 \times 0.448 - 0.211 \times 0.01495 = -0.03396 \text{ kJ kg}^{-1} \text{ K}^{-1}. \end{aligned}$$

Solution

Since the process is adiabatic and reversible we will use the relation

$$\begin{split} &(\partial T/\partial P)_s = Tv\beta_P/c_P \text{ or } \\ &(\Delta T/\Delta P)_s = T_o v\beta_P/c_P. = \\ &\{250K\times1.1\times10^{-4}m^3kg^{-1}\times48\times10^{-6}K^{-1} \ \} \ /0.372kJkg^{-1}K^{-1} = 3.548\times10^{-6} \ K/Kpa \ and \\ &\Delta P = 100\times1000 \ Kpa. \ Hence \\ &dT_s = 0.36 \ K \ and \ T \ will \ rise \\ &to 250.36 \ K. \end{split}$$

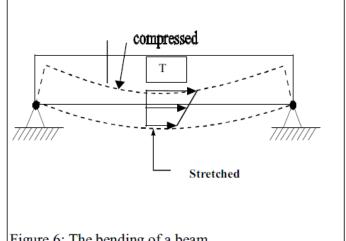


Figure 6: The bending of a beam.

Applying the First law to an adiabatic reversible process

 $du_s = -Pdv_s$ .

Recall from Eq. (37) that

 $dv_s = -dT_s \beta_T c_v/(T \beta_P)$ , i.e.,

 $du_s = \{P \beta_T c_v / (T \beta_P)\} dT_s$ 

Integrating and assuming that P is not a function of temperature and remaining at an average value of 500 bar.

 $du_s = \{500 \text{ bar} \times 7.62 \times 10^{-7} \text{ bar}^{-1} \times 0.364 \text{ kJ kg}^{-1} \text{ K}^{-1} \}$ 

 $\div (250 \text{ K} \times 48 \times 10^{-6} \text{ K}^{-1})$ } 0.36 K = 0.00416 kJ kg<sup>-1</sup>.

The temperature after the load is removed is 250 K, since the process is reversible.